Aerodynamic aspects of the sealing of gasturbine rotor-stator systems Part 3: The effect of nonaxisymmetric external flow

on seal performance

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Received 3 April 1987 and accepted for publication on 29 June 1987

The effect of a nonaxisymmetric external flow on the sealing performance of a shrouded rotor-stator system has been studied using flow visualization and pressure measurements. When ingress occurred, ingested fluid moved transversely across the wheel space from regions of high to low pressure in the external-flow annulus. One of the four seal geometries tested contained a double-shrouded seal, and the inner shroud attenuated the external pressure asymmetries so that most of the ingested fluid was confined to the space between the shrouds. Six different circumferential variations of pressure were used to separate the effects of flow rate and pressure asymmetry, and it was found that, for all four seals, the minimum sealing-flow rate necessary to prevent ingress could be correlated with the maximum circumferential pressure difference measured in the external-flow annulus.

Keywords: sealing; rotating disks; rotor-stators

Introduction

Studies of ingress in a number of different shrouded rotor-stator systems were reported in Parts 1 and 2 (these parts, Refs. 1 and 2, are referred to here as I and II) for the case where the external flow was either zero or quasi-axisymmetric. It was shown that at small values of $\text{Re}_{\omega}/\text{Re}_{\theta}$ (where Re_{ω} and Re_{θ} are the externalflow and rotational Reynolds numbers, respectively) there is rotation-dominated regime where $C_{w,min}$, the minimum sealingflow rate necessary to prevent ingress, increases with increasing Re_{θ} . At large values of Re_{w}/Re_{θ} , there is an external-flowdominated regime in which $C_{w,min}$ is proportional to Re_w but is virtually independent of $Re₉$.

For the tests reported in II, perfect axisymmetry of the external flow was not achieved. As the small variations in circumferential pressure increased with increasing Re_w, it was not possible to conclude whether the results for the externalflow-dominated regime were caused by Re_w or by the pressure asymmetries. It is the object of the work reported below to determine which of these effects causes ingress.

Of the seven geometries tested in I, seals 1, 2, 3, and 5 were examined in the present study, and these geometries are illustrated in Figure 1: see Refs. 1, 2, or 3 for details of the experimental apparatus. The second and third sections below, respectively, describe the pressure asymmetries and their effect on $C_{w,min}$.

Pressure asymmetries

Generation of the asymmetries

A contraetion with a 16:1 area ratio was fitted to the upstream end of the external flow annulus, as shown in Figure 2, and the inlet of the contraction was covered with wire mesh. By attaching sections of "honeycomb" to parts of the mesh, a number of different pressure distributions could be generated in the circumferential direction.

Figure 3 shows the circumferential variation of static pressure, measured at the 27 taps on the outside wall of the annulus, for six different patterns of pressure asymmetry. The measurements were obtained for seal 1 (the datum seal) with $G_c = 0.01$, Re_g= 0, and Re_w = 1.1 × 10⁶. Only static pressures were measured, and no attempt was made to measure the variation of total pressure or secondary flow in the annulus. For further details of the flows in the annuli of turbomachinery, see Refs. 4 and 5.

The spectral distribution is different for each of the six pressure asymmetries, and one simple measure of the asymmetry is Δp_{max} , the maximum pressure difference between any two of the 27 taps. To give some idea of the asymmetry of the pressure distributions, it is convenient to define a pressure coefficient

$$
C_{p,\text{max}} = \Delta p_{\text{max}} / \frac{1}{2} \rho \bar{W}^2 \tag{1}
$$

The asymmetries created by the honeycomb sections are numbered 1 to 6 in order of increasing $C_{p,\text{max}}$, and Table 1 shows the values obtained for seal 1 with $G_c=0.01$, Re_{$\theta=0$}, and $Re_w = 1.1 \times 10^6$. For the quasi-axisymmetric external flow discussed in II, the pressure asymmetries were caused mainly by small misalignments in the external flow passages and not by the deliberate introduction of disturbances into the flow. The quasiaxisymmetric case, referred to as asymmetry number 0, is also included in Table 1. It was found that $C_{p,\text{max}}$ was sensibly independent of Re_{w} and, for large values of Re_{w} , it was not significantly affected by the speed of the rotor.

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Figure I Simplified diagram of the four seals tested: (a) seal 1 ; (b) seal 2; (c) seal 3; (d) seal 5

Figure 2 External-flow rig used for the asymmetry tests

The effect of the asymmetries on ingress

As reported in II, flow visualization revealed that, with an external flow, ingress could occur even when the rotor was stationary. Smoke introduced into the external flow showed that fluid moved transversely across the wheel space from regions of high to low pressure in the annulus. In the tests

Figure 3 The circumferential variation of static pressure in the annulus for each of the six asymmetries: $C_{\mathsf{w}}{=}0$, $\mathsf{Re}_\theta{=}0$, Re_w=1.1 × 10⁶ Symbol $x \in \Box \quad \diamond \quad \rhd + \quad \prec$
Asymmetry number $x = 1$ and $x = 2$ and $x = 3$ and $x = 6$ Asymmetry number

reported here, different asymmetries caused different effects with respect to the extent and direction of the ingested flow.

Figure 4 shows typical examples of the relationship between the observed flow patterns and the circumferential pressure distribution for seal 1 with $G_c=0.01$, $C_w=0$, $Re_\theta=0$, and $Re_w = 1.1 \times 10^6$ for asymmetry numbers 3 and 5. It is apparent

Table 1 The value of $C_{p,\text{max}}$ for each of the asymmetries for seal 1 with $G_c = 0.01$, Re₀=0, and Re_w=1.1 × 10⁶

Figure4 Observed flow patterns in the wheel space and circumferential variation of static pressure in the annulus for $C_w=0$, Re_θ =0, Re_w =1.1 × 10⁶: (a) asymmetry no. 3; (b) asymmetry no. 5

that the flow direction in both cases is from the high- to the lowpressure region in the annulus: the wheel space acts as a short circuit to redress the balance of pressure in the annulus.

The above examples were for the case where $C_w=0$. If the cooling flow rate is increased, the amount of ingested fluid can be reduced and eventually ingress is prevented. Figure 5 shows sketches of the flow patterns deduced from flow visualization for seal 1 with $G_c = 0.01$ and for seal 2, the double-shrouded seal, with $G_{c1} = G_{c2} = 0.01$ for $\text{Re}_\theta = 0$ and $\text{Re}_w = 0.21 \times 10^6$. Referring to the diagrams for seal 1, note that at $C_w \approx 900$ the amount of ingestion is reduced and at $C_w \approx 1800$ it is virtually eliminated.

The patterns for seal 2 are different to those for seal 1: the outer wheel space provides a "ring road" that allows the transfer of most of the flow from the high- to the low-pressure regions in the external-flow annulus. In fact, at $C_w \approx 900$ no external flow enters the inner wheel space, and at C_w -1800 all ingress has ceased. In the absence of external flow, it was shown in I that seal 2 had no particular merits compared with other seals. However, with external flow, seal 2 has the property of reducing ingress into the inner wheel space, and this could be used to obvious advantage in gas-turbine cooling systems.

From tests conducted for seal 2, it was found that pressure asymmetries measured on the stator between the two shrouds were much smaller than those measured on the outer wall of the annulus. As a consequence, the pressure asymmetry experienced by the fluid in the inner wheel space is significantly attenuated. As the magnitude of the pressure asymmetry affects the amount of ingestion that occurs, the inner wheel space is "buffered" by the second shroud.

Figure 6 shows the effect of C_w on the circumferential variation of Δp_1 , where Δp_1 is the local pressure difference measured between a tap on the stator at $x=0.97$ and the corresponding tap on the outer wall of the annulus. The measurements were made for seal 1 with $G_c=0.01$, Re $_0=0$, and $Re_w = 1.1 \times 10^6$, and the results for asymmetries 1, 3, and 5 are shown. It can be seen that, for all three cases, Δp_1 is negative at all circumferential locations for $C_w=0$. Increasing C_w causes an increase in the magnitude of Δp_1 , but the circumferential distribution is virtually unaltered.

The distribution of Δp_1 is related to the ingestion observed by flow visualization, and the determination of $C_{w,\text{min}}$ is described below.

Figure 5 The effect of Cw on the observed flow patterns in the wheel space for $G_{\rm c}$ =0.01, Re $_{\theta}$ =0, Re $_{\rm w}$ =0.2×10°: (a) $C_{\rm w}$ =0; (b) $C_w \approx 900$; (c) $C_w \approx 1800$

Figure 6 The effect of Cw on the circumferential variation **of** pressure drop across seal 1 for G_c =0.01, Re $_\theta$ =0, Re_w=1.1 × 10⁶: (a) asymmetry no. 1; (b) asymmetry no. 3; (c) asymmetry no. 5

1 2 3 4 5 6

Figure 7 The variation of $C_{\mathsf{w}, \mathsf{min}}$ with Re_w for the different **asymmetries:** $G_c = 0.01$, $Re_\theta = 0$
Symbol Symbol $x \cup \diamond$ \rightarrow $+$ \lhd

Figure 8 The effect of Re_{θ} on the variation of C_{w, min} with Re_w for **seal 1** with asymmetry number **6:**

The determination of $C_{w,min}$

Flow visualization

Asymmetry number 0

The flow visualization technique was the same as that described in I and II: smoke was injected into the external flow, and $C_{w,min}$ was taken to be the minimum value of C_w that prevented the appearance of smoke in the illuminated wheel space.

Figure 7 shows the variation of $C_{\mathbf{w},\mathbf{min}}$ with Re_w for the quasiaxisyrnmetric case (number 0) and for each of the six asymmetries tested. The results, which were Obtained for seal 1 with $G_c = 0.01$ and $\text{Re}_\theta = 0$, show that $C_{w,\text{min}}$ is proportional to Re_w and increases monotonically from asymmetry number 0 to number 6: that is, $C_{w,min}$ increases with increasing $C_{p,max}$ as well as with increasing Re_{w} .

Figure 8 shows the effect of Re_{θ} on the variation of $C_{\text{w,min}}$ with Re_w, for seal 1 with $G_c = 0.01$, for asymmetry number 6. For comparison, a curve representing the average results for the quasi-axisymmetric case, with $Re_\theta=0$, is shown. For small values of Re_w/Re_θ the results show the sealing effect of the external flow: $C_{\mathbf{w},\text{min}}$ decreases with increasing $\mathbf{Re}_{\mathbf{w}}$. For large values of Re_w/Re_θ , in the external-flow-dominated regime, the effect of Re_{θ} is small and $C_{w,\text{min}}$ becomes proportional to Re_w. However, as discussed above, the value of $C_{w,min}$ required for asymmetry number 6 is significantly greater than that for the quasi-axisymmetric case.

The effect of Re_{θ} on the variation of $C_{w,\text{min}}$ with Re_w for the double-shrouded seal 2 with asymmetry number 4 is shown in Figure 9 for the case of $G_{c1} = G_{c2} = 0.01$. In Figure 9(a), for the

outer shroud, Re₀ and $C_{w,min}$ are based on the outer radius r_o ; in Figure 9(b), for the inner shroud, Re_{θ}^{*} and $C_{w,\text{min}}^{*}$ are based on the radius r_o^* of the inner shroud ($r_o^* = 0.9r_o$). For $Re_w = 0$, there is little difference in performance between the two shrouds: $C_{\text{w,min}}$ and $C_{\text{w,min}}^*$ are proportional to Re_e and Re_e^{*}, respectively. For large values of Re_{w} , the value of $C_{w,min}$ for the outer shroud is much greater than that for the inner shroud. This is consistent with the results discussed in the previous section in which flow visualization revealed that the outer wheel space acted as a ring road that transported most of the ingested fluid.

Pressure measurements

Flow visualization showed that $C_{w,min}$ increased with increasing $C_{p,\text{max}}$, and a simple model for this behavior is presented below.

It is assumed that, in the external-flow-dominated regime, the circumferentially averaged pressure drop Δp_s across the shroud is related to the volumetric flow rate Q of the sealing air by

$$
Q = C_d A (2\Delta p_s/\rho)^{1/2} \tag{2}
$$

where $A = 2\Pi r_0 s_c$ is the area of the seal opening and C_d is a discharge coefficient. It follows from Equation 2 that

$$
C_{w} = 2\sqrt{2} \Pi C_{d} G_{c} (\rho \Delta p_{s} r_{o}^{2}/\mu^{2})^{1/2}
$$
 (3)

and it is postulated that $C_w = C_{w,\text{min}}$ when $\Delta p_s = C \Delta p_{\text{max}}$, where C is an empirical constant. Hence

$$
C_{\rm w,min} = 2\Pi K G_{\rm c} P_{\rm max}^{1/2} \tag{4}
$$

where

$$
K = \sqrt{2C} C_{\rm d} \tag{5}
$$

Figure 9 The effect of Re₀ on the variation of $C_{\mathbf{w},\text{min}}$ with Re_w for seal 2 with asymmetry number 0 for $G_c = 0.01$; (a) outer shroud; (b) **inner** shroud Symbol $\hspace{.15cm} \circ \hspace{.15cm} \square \hspace{.15cm} \triangledown \hspace{.15cm} \triangle \hspace{.15cm} \times \hspace{.15cm} \blacktriangledown$ \blacktriangle **Symbol** $\begin{array}{ccccccc}\n\bigcirc & \Box & \bigtriangledown & \bigtriangleup & \times & \blacktriangledown & \blacktriangle & \mathsf{R}\n\mathsf{e}_{\theta}/10^6 & & & 0.2 & 0.4 & 0.6 & 0.8 & 1.0 & 1.2\n\end{array}$

and

$$
P_{\text{max}} = \rho \Delta p_{\text{max}} r_o^2 / \mu^2 = \frac{1}{2} C_{p,\text{max}} \text{Re}_w^2
$$
 (6)

Figure 10 shows the variation of $C_{w,\text{min}}$ with the parameter $2\pi G_c P_{\text{max}}^{1/2}$. The measured values of $C_{\text{w,min}}$ were obtained by flow visualization for seals 1,2, 3 and 5 (for seal 2, only the results for the outer wheel space are included) with G_c =0.005, 0.01, and 0.02, $Re_{\theta} = 0$ and $Re_{w} \le 1.1 \times 10^{6}$. Despite the scatter, the results are correlated reasonably well using a value of $K = 0.6$ in Equation 4. Since this equation takes no account of the seal geometry or of the spatial distribution of the pressure asymmetry, the results suggest that these effects may be of secondary importance and that P_{max} can be used to estimate $C_{\text{w,min}}$.

More tests need to be carried out with systematic variation of pressure and swirl in the external-flow annulus, but Equation 4 may prove to be helpful for design purposes in the external-flowdominated regime. In the rotation-dominated regime, correlations of the form

$$
C_{\rm w,min} \propto G_{\rm c}^m \text{Re}_\theta^n \tag{7}
$$

were determined for the six different seal geometries in I. A conservative design would be one in which $C_{w,\min}$ was assumed to be the greater of the values determined for the two regimes.

Conclusions

Six different circumferential distributions of static pressure have been used to separate the effects of flow rate and pressure asymmetry in the external-flow annulus of a shrouded rotatingdisk system. Flow visualization revealed that, after ingress had occurred, ingested fluid moved transversely across the wheel space from regions of high to low pressure in the annulus. For seal 2, the double-shrouded seal, the outer wheel space between the two shrouds acted as a ring road allowing most of the ingested fluid to be transferred from regions of high to low pressure without crossing the inner wheel space.

In the external-flow-dominated regime, it was found that $C_{\text{w,min}}$ increased with increasing P_{max} , where P_{max} is the dimensionless maximum circumferential static pressure difference measured in the external-flow annulus. A simple correlation between $C_{w,min}$, G_c , and P_{max} was obtained for the four seal geometries tested with $G_c=0.005$, 0.01, and 0.02 and $Re_w < 1.1 \times 10^6$. Using this correlation, together with correlations for the rotation-dominated regime studied in I, it should be possible to obtain a conservative estimate of $C_{\mathbf{w},\text{min}}$. From the results of the tests reported in this series of papers, it would appear that a double-shrouded radial clearance seal should have advantages over the other seal geometries in minimizing the amount of hot gas that is ingested into the wheel spaces of gas turbines.

It should be emphasized, however, that the above tests were conducted with spatially aperiodic pressure distributions. In a gas turbine, nozzles, and combustion chambers often create pressure disturbances that contain swirl and spatial periodicity. Attempts to model such disturbances, theoretically and experimentally, and to determine their effect on the "ingress problem" are continuing in the Thermo-Fluid Mechanics Research Centre at the University of Sussex.

Acknowledgments

The authors wish to thank Rolls Royce plc, Ruston Gas Turbines pie, and the Science and Engineering Research Council for sponsoring the work described in this series of papers. They are also indebted to Dr. J.R. Pincombe for his assistance with the flow visualization and to the technicians in the School of Engineering and Applied Sciences who manufactured the experimental apparatus.

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